



Heat Generation/Absorption Effect on Natural Convection Heat Transfer in a Square Enclosure Filled with a Ethylene Glycol - Copper Nanofluid Under Magnetic Field

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Abstract: This paper examines the natural convection in a square enclosure that is filled with a nanofluid. This nanofluid with Ethylene Glycol based containing Copper nanoparticle is influenced by a uniform horizontal magnetic field and uniform heat generation or heat absorption. The enclosure is bounded by two isothermal vertical walls at different temperatures and by two horizontal adiabatic walls. The governing equations needed to deal this problem (mass, momentum, and energy) are solved numerically using the commercial simulation software COMSOL Multiphysics. In order to increase the natural convective heat transfer in a square cavity, the effect of heat generation or absorption on the isothermal, streamline contours and the Nusselt number are studied when the Prandtl number is $Pr = 151$.

Keywords: Heat Transfer, Natural Convective, Square Enclosure, EG-Cu Nanofluid, Magnetic Field, Generation/Absorption, Comsol Multiphysics

1. Introduction

The study of magnetohydrodynamic flow when especially associated with heat transfer has attracted several researchers in recent years due to its wide variety of applications in engineering areas and technology such as plasma industries, cooling of nuclear reactor, optimization of crystal growth processes, electronic package, geothermal energy extractions, solidification of metal alloys, metallurgical applications involving casting and solar technology.

The existence of an external magnetic field is used as a control mechanism in material manufacturing industry, as the convection currents are suppressed by Lorentz force which is produced by the magnetic field.

Rudraiah et al. [1] studied numerically the natural convection of an electrically conducting fluid in the presence of a magnetic field parallel to gravity. They have pointed out that the average Nusselt number decreases with an increase in the Hartmann, and the Nusselt number approaches one under a strong magnetic field. Piazza and Ciofalo [2]

considered the buoyancy-driven magnetohydrodynamic flow in a liquid-metal-filled cubic enclosure differentially heated at two opposing vertical walls. Kandaswamy et al. [3] studied the magnetoconvection flow in a enclosure with partially active vertical walls. They have considered nine different positions of active zone for different value of the Rayleigh and Hartman numbers. It has been found that the average Nusselt number decreases with an increase of Hartman number and increases with the Rayleigh number. Also, the convection mode of heat transfer is converted into a conduction mode for an enough large magnetic field. Ghassemi et al. [4] investigated the effect of the magnetic field on natural convection in a enclosure with two adiabatic baffles. They showed that the magnetic field had an adverse effect on the Nusselt number, particularly at low values of Rayleigh number. Pirmohammadi and Ghassemi [5] considered the effect of the magnetic field on convection heat transfer inside a tilted square enclosure. Their study showed

that the heat transfer mechanism and flow characteristics inside the enclosure depend strongly upon both magnetic field and inclination angle. Sathiyamoorthy and Chamkha [6] investigated natural magneto-convection of liquid Gallium in a square heated cavity. They found that the average Nusselt number decreased nonlinearly with the Hartman number. Sivasankaran and Ho [7] studied numerically the effects of temperature dependent properties on the natural convection of water in a enclosure under the influence of a magnetic field. They showed that the heat transfer rate decreases with an increase of the magnetic field and influenced by the direction of the external magnetic field. Sarris *et al* [8] investigated the natural convection of an electrically conducting fluid in a laterally and volumetrically heated square cavity under the influence of a magnetic field. They concluded that the heat transfer is enhanced with increasing internal heat generation parameter, but no significant effect of the magnetic field is observed due to the small range of the Hartmann numbers. Bhuvanewari *et al.* [9] investigated magnetic convection in an enclosure with non--uniform heating on both walls. They found that the heat transfer rate is increased on increasing the amplitude ratio. Kolsi et al [10] studied the effect of magnetic field orthogonal to the isothermal walls on the unsteady natural convection in a differentially heated cubic cavity. They observed the role of magnetic field on the damping and laminarization of flow field.

The main fluids used for heat transfer application are water, ethylene glycol, mineral oil or propylene glycol have a low thermal conductivity. In order to improve the heat transfer performance of inherently poor conventional heat transfer base fluids, the use of nanofluids is one of the most effective mechanisms which has a superior thermal conductivity compared to the base fluid.

Many publications, either theoretical [11,12] or experimental [13-16], use the water as base fluid with several particles such as Al₂O₃ [17, 18], Cu [19], CuO [20], TiO₂ and Ag. But, a few research work on ethylene glycol [21-26] as base fluid despite its extensive application in industrial and energy saving perspectives such as in power generation, heating and cooling processes, heat exchanger, chemical processes, electronics, transportation, automotive and others.

In addition, the influence of a heat generation in a cavity containing only base fluid without nanofluids was studied by a few researchers such as Khanafer *et al.* [27] and Hossain *et al.* [28].

Therefore, The main originality of the present study is to investigate in detail the natural convection in a square enclosure by utilizing Ethylene Glycol - Copper Nanofluid

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = \frac{1}{\rho_{nf}} \left[-\frac{\partial p}{\partial y} + \mu_{nf} \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right) + (\rho\beta)_{nf} g(T - T_c) - \sigma_{nf} B_0^2 v \right] \quad (3)$$

Where T_c is the temperature of cold wall , β is the thermal expansion coefficient, σ_{nf} is the electrical conductivity of nanofluid and B₀ is the magnetic field

and is influenced by a horizontally applied uniform magnetic field in the presence of uniform heat generation or absorption. The physical models are solved using the commercial simulation software COMSOL Multiphysics. The results of the simulation are presented and discussed in the sections below.

2. Numerical Model

2.1. Simplifying Assumptions

The governing equations of the numerical model are written taking into account the following simplifying assumptions:

- The base fluid Ethylene Glycol and the nanoparticles are assumed in thermal equilibrium.
- The nanofluid is Newtonian and incompressible.
- The flow is considered to be steady, two-dimensional and laminar.
- The radiation effects are negligible.
- The displacement currents, induced magnetic field, dissipation and joule heating are neglected.

2.2. Governing Equations

2.2.1. Equations of the Model in Dimensional Form

The governing equations for this problem are based on the balance laws of mass, momentum and thermal energy of a steady laminar flow. Taking into account the assumptions previously-mentioned, the simplified system of equations is written in two-dimensional form as follows:

Mass conservation equation:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

Where u and $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$ v are the velocities components in the x and the y direction, respectively.

Momentum conservation equation in the x direction :

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \frac{1}{\rho_{nf}} \left[-\frac{\partial p}{\partial x} + \mu_{nf} \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right) \right] \quad (2)$$

Where ρ_{nf} , p and μ_{nf} are the density of nanofluid, fluid pressure and the effective dynamic of the nanofluid, respectively.

Momentum conservation equation in the y direction :

strength.

Energy conservation equation:

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \alpha_{nf} \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + \frac{q(T - T_c)}{(\rho C_p)_{nf}} \quad (4)$$

$$\mu_{nf} = \frac{\mu_f}{(1 - \phi)^{2.5}} \quad (11)$$

Where α_{nf} is the thermal diffusivity of nanofluid and q is the heat generation or the heat absorption.

The properties of the nanofluid used in this study are inspired by classical models reported in the literature [29].

$$\sigma_{nf} = (1 - \phi) \sigma_f + \phi \sigma_p \quad (5)$$

$$\rho_{nf} = (1 - \phi) \rho_f + \phi \rho_p \quad (6)$$

$$(\rho C_p)_{nf} = (1 - \phi) (\rho C_p)_f + \phi (\rho C_p)_p \quad (7)$$

$$(\rho \beta)_{nf} = (1 - \phi) (\rho \beta)_f + \phi (\rho \beta)_p \quad (8)$$

$$\alpha_{nf} = \frac{k_{nf}}{(\rho C_p)_{nf}} \quad (9)$$

In the above equations, k is the thermal conductivity, C_p is the specific heat and ϕ is the solid volume fraction. The subscripts f and p means the fluid and nanoparticle, respectively.

The thermal conductivity and the effective dynamic viscosity of the nanofluid can be modelled by [30, 31].

$$k_{nf} = k_f \left[\frac{(k_p + 2k_f) - 2\phi(k_f - k_p)}{(k_p + 2k_f) + \phi(k_f - k_p)} \right] \quad (10)$$

$$u^* \frac{\partial v^*}{\partial x^*} + v^* \frac{\partial v^*}{\partial y^*} = -\frac{\partial p^*}{\partial y^*} + \frac{\mu_{nf}}{\alpha_{nf} \alpha_f} \left(\frac{\partial^2 u^*}{\partial x^{*2}} + \frac{\partial^2 u^*}{\partial y^{*2}} \right) + \frac{(\rho\beta)_{nf}}{\rho_{nf} \beta_f} Ra Pr T^* - Ha^2 Pr v^* \quad (14)$$

Energy conservation equation:

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{\alpha_{nf}}{\alpha_f} \left(\frac{\partial^2 T^*}{\partial x^{*2}} + \frac{\partial^2 T^*}{\partial y^{*2}} \right) + \frac{\alpha_{nf}}{\alpha_f} T^* q^* \quad (15)$$

$$u^* = \frac{\partial \psi}{\partial y^*} \text{ et } v^* = -\frac{\partial \psi}{\partial x^*} \quad (18)$$

These previous equations are solved using the commercial simulation software COMSOL MULTIPHYSICS.

2.2.3. Nusselt Number Calculations

The local Nusselt number on the left hot wall can be defined as :

$$Nu_y(y^*) = -\frac{k_{nf}}{k_f} \left(\frac{\partial T^*}{\partial x^*} \right)_{x^*=0} \quad (16)$$

The average Nusselt number (Nu_m) is calculated by integrating Nu_y along the hot wall :

$$Nu_m = \int_0^1 Nu_y(y^*) dy^* \quad (17)$$

The stream function is determined from :

2.2.2. Equations of the Model in Non-Dimensional Form

The governing equations are nondimensionalized using the following variables:

$$x^* = \frac{x}{L}, y^* = \frac{y}{L}, u^* = \frac{uL}{\alpha_f}, v^* = \frac{vL}{\alpha_f}, T^* = \frac{T - T_c}{T_h - T_c},$$

$$p^* = \frac{pL^2}{\rho_{nf} \alpha_f^2}, Pr = \frac{\nu_f}{\alpha_f}, Ha = B_0 L \sqrt{\frac{\sigma_{nf}}{\rho_{nf} \nu_f}},$$

$$Ra = \frac{g \beta_f L^3 (T_h - T_c)}{\nu_f \alpha_f} \text{ and } q^* = \frac{qL^2}{\alpha_{nf} (\rho C_p)_{nf}}$$

Based on the presumptions above, the dimensionless governing equations can be expressed as follows:

Mass conservation equation:

$$\frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} = 0 \quad (12)$$

Momentum conservation equation in the x direction :

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = -\frac{\partial p^*}{\partial x^*} + \frac{\mu_{nf}}{\alpha_{nf} \alpha_f} \left(\frac{\partial^2 u^*}{\partial x^{*2}} + \frac{\partial^2 u^*}{\partial y^{*2}} \right) \quad (13)$$

Momentum conservation equation in the y direction :

2.3. Boundary Conditions of the Model

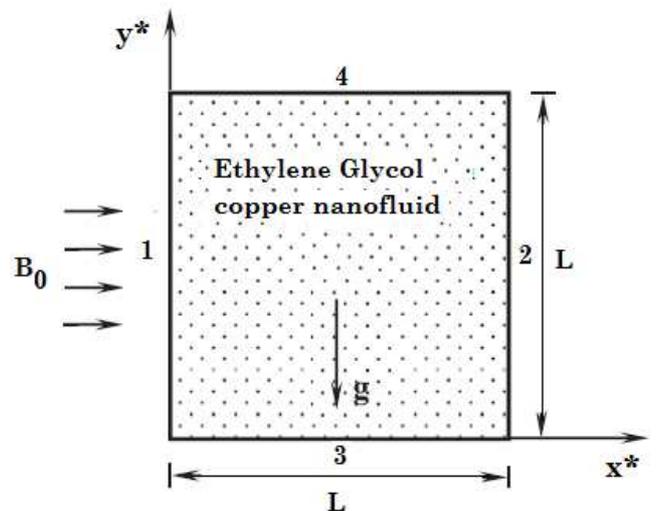


Figure 1. Schematic diagram of the physical model.

The schematic of a two-dimensional square cavity is shown in Figure 1. It is filled with an electrically conducting fluid and a Ethylene Glycol - Copper nanofluid. The left and right side walls are kept at constant temperatures, T_C and T_h , respectively. The top and bottom surfaces are assumed to be

insulated and impermeable. A magnetic field with uniform strength B_0 is applied in the x direction which is perpendicular to the gravity.

The boundary conditions used for dimensional and dimensionless form are summarized to the Table 1.

Table 1. Boundary conditions for dimensional and dimensionless form.

Border	Dimensional form			Dimensionless form		
	Condition on u	Condition on v	Condition on T	Condition on u^*	Condition on v^*	Condition on T^*
1	0	0	T_h	0	0	1
2	0	0	T_C	0	0	0
3	0	0	$\frac{\partial T}{\partial y}$	0	0	$\frac{\partial T^*}{\partial y^*}$
4	0	0	$\frac{\partial T}{\partial y}$	0	0	$\frac{\partial T^*}{\partial y^*}$

The thermo-physical properties of base fluid (Ethylene Glycol: $C_2H_6O_2$) [32, 25] with Copper [33] (Cu), Alumina [34] (Al_2O_3), Titanium oxide (TiO_2) and Copper oxide (CuO) [35] are given in Table 2.

Table 2. Thermo-physical properties of Ethylene Glycol and nanoparticles.

		Pr	ρ [kg/m ³]	Cp [J/kg.K]	K [W/m.K]	$\beta \times 10^{-5}$ [K ⁻¹]	$\alpha \times 10^{-5}$ [m ² /s]
Base fluid	Ethylene Glycol ($C_2H_6O_2$)	151	1109	2400	0.26	65	0.0214
	Copper (Cu)		8933	385	401	1.67	11.7
Nano-particles	Alumina (Al_2O_3)		3970	765	40	0.85	1.3
	Titanium oxide (TiO_2)		4250	686.2	8.9538	0.9	0.31
	Copper oxide (CuO)		6500	535.6	20	0.85	0.57

3. Validation of the Model

In order to validate the present numerical method for natural convection in an enclosure, comparisons with the experimental results from Linthorst et al. [36] are made. They measured the dimensionless y -direction velocity by means of a laser doppler velocimeter at midplane ($y^* = 0.5$) of a enclosure filled with air. For the computational validation, the Rayleigh numbers are set to 1.3×10^5 . A square cavity with a length $L = 4$ cm, whose upper and bottom walls are adiabatics, is chosen. The Prandtl number of the air is fixed as 0.733 while the thermal diffusivity is 26.89×10^{-6} m²/s, which was deduced from the mean temperature of fluid equal to 320 K. The computational domain is divided by 51×51 uniform grids spatially and the same grids are used for further investigation.

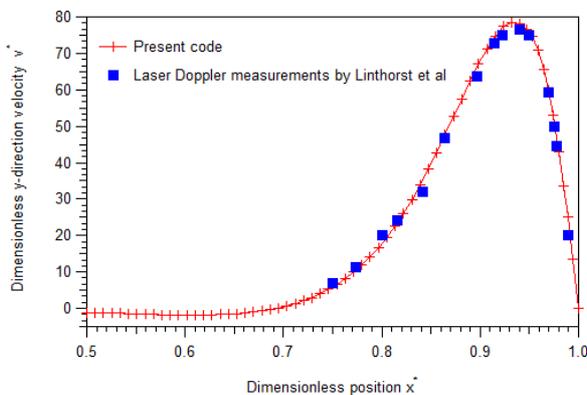


Figure 2. Comparison of dimensionless y -direction velocities at midplane of a enclosure for Rayleigh number $Ra = 1.3 \times 10^5$.

The figure 2 shows the dimensionless y -direction velocity according to dimensionless position x^* at midplane of a cavity for Rayleigh number $Ra = 1.3 \times 10^5$ between the theoretical values of our code and the values measured by Linthorst et al. [36]. From this figure, we notice the quite satisfactory agreement between the calculated and measured values, thereby justifying the different assumptions adopted.

4. Results and Discussion

In this study, the Prandtl number is assumed to be $Pr = 151$. The Rayleigh number (Ra), the Hartmann number (Ha), the dimensionless heat generation or absorption and the solid volume fraction (ϕ) are assumed to be in the following ranges: $10^3 \leq Ra \leq 10^7$, $0 \leq Ha \leq 80$, $-12 \leq q^* \leq 12$ and $0 \leq \phi \leq 0.06$. The effect of Hartmann number and Rayleigh number on the isothermal and streamline contours, effect of magnetic field on the Nusselt number, effect of heat generation or absorption and Rayleigh number on the isothermal and streamline contours, effect of heat generation or absorption and Hartmann number on the Nusselt number and the comparison with other nanofluids are presented and discussed.

4.1. Effect of Hartmann Number and Rayleigh Number on the Isothermal and Streamline Contours

Figure 3 shows the isotherms for different Hartmann number ($Ha = 0, 20, 40, 60$ and 80) and Rayleigh number ($Ra = 10^3, 10^4, 10^5, 10^6$ and 10^7) at dimensionless heat generation $q^* = 2$, a solid volume fraction $\phi = 0.05$ and Prandtl number $Pr = 151$.

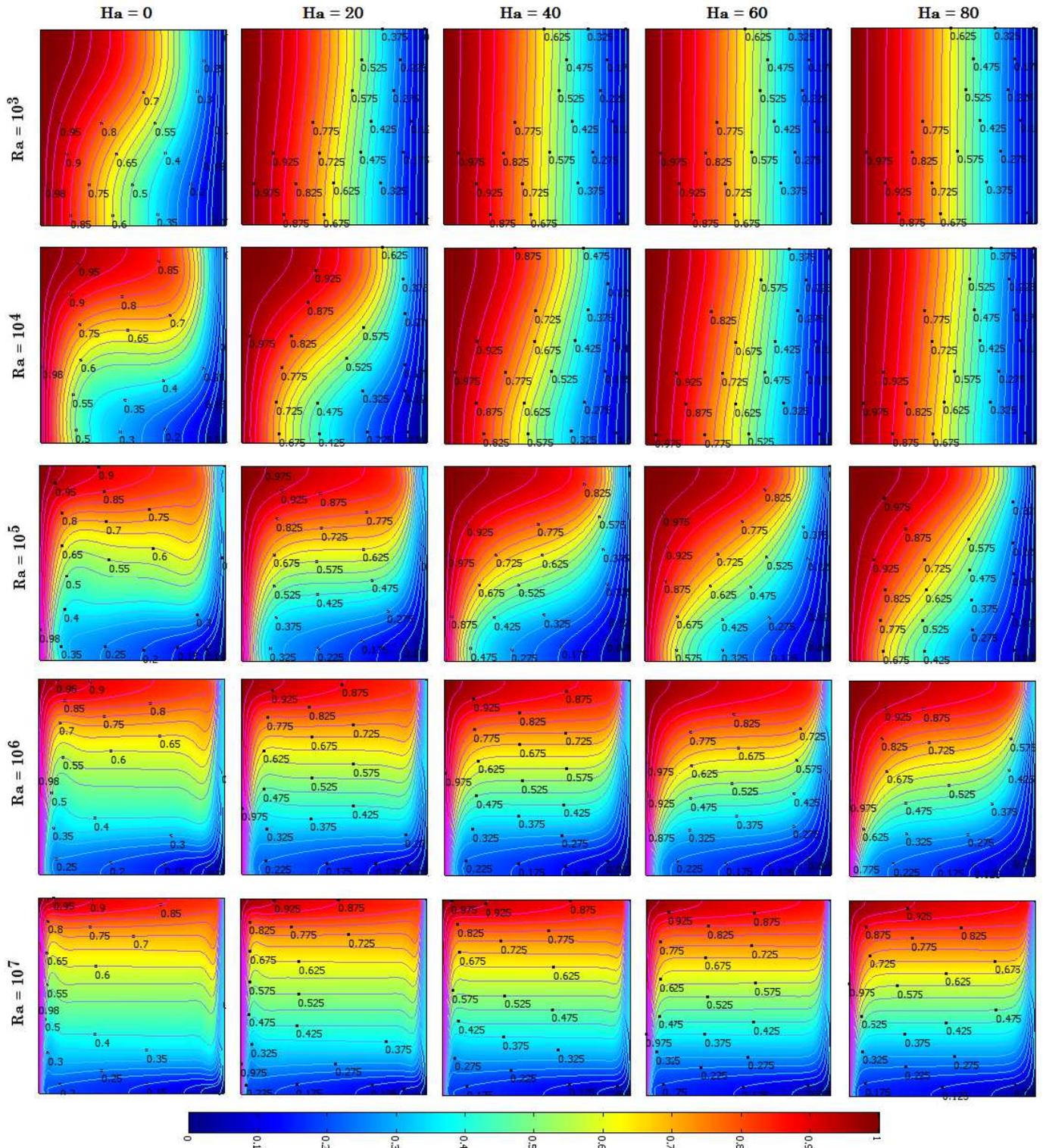


Figure 3. Isotherms for different Hartmann number and Rayleigh number at dimensionless heat generation $q^* = 2$ and a solid volume fraction $\Phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

Figure 4 shows the streamlines for different Hartmann number ($Ha = 0, 20, 40, 60$ and 80) and Rayleigh number ($Ra = 10^3, 10^4, 10^5, 10^6$ and 10^7) at dimensionless heat generation $q^* = 2$, a solid volume fraction $\Phi = 0.05$ and Prandtl number $Pr = 151$.

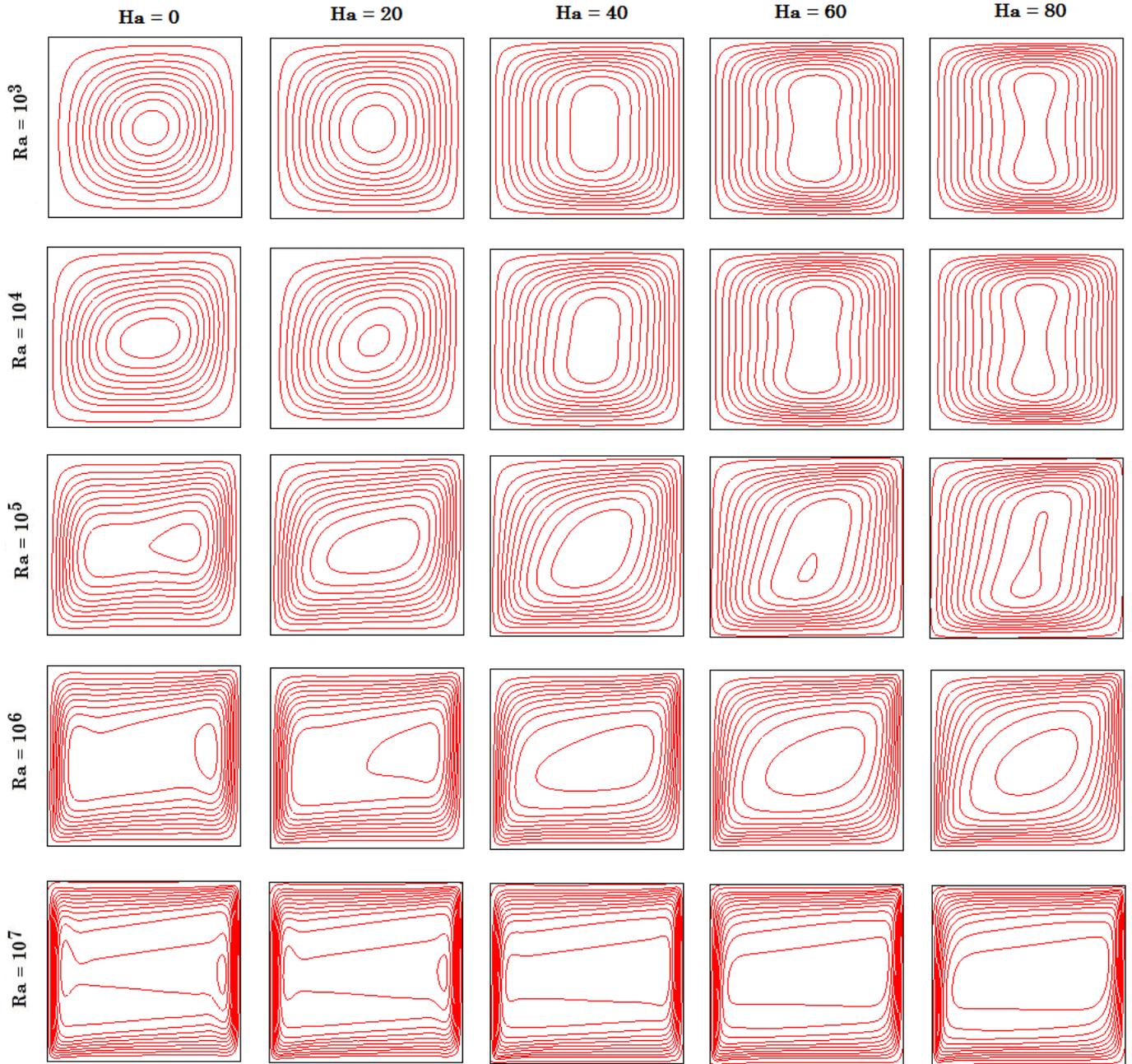


Figure 4. Streamlines for different Hartmann number and Rayleigh number at dimensionless heat generation $q^* = 2$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From these figures, we observe that there is conduction dominated regime at low Rayleigh numbers and a convection dominated regime at high Rayleigh numbers. Indeed, for low Rayleigh number $Ra = 10^3$, the isotherms are parallel vertical lines with a weak clockwise circulation in the cavity. Also, we watch that the circulations strength decreases with the increasing of the Hartmann number and increases when the Rayleigh number increases. In addition, the sign of Hartmann number is opposite to the sign of Rayleigh number in source term of equation (14). So, there is an opposite effect of these two parameters on flow regime and Nusselt number. At Rayleigh number $Ra = 10^5$, the convection mode is pronounced, the isotherms change from

horizontal to vertical when the Hartmann numbers increase. Also, the flow cell becomes stronger. At higher Rayleigh number 10^7 , the convection is dominant and the circulating cell becomes very strong. Also, the streamlines are crowded near the enclosure wall and the core is empty. As well as isotherms are stratified in vertical direction except near the insulated surfaces of the cavity and appear as horizontal lines in the enclosure core.

4.2. Effect of Magnetic Field on the Nusselt Number

Figure 5 shows the Average Nusselt number according to the Rayleigh number ($Ra = 10^3, 10^4, 10^5, 10^6$ and 10^7) for different Hartmann number ($Ha = 0, 20, 40, 60$ and 80) at

dimensionless heat generation $q^* = 2$, a solid volume fraction $\phi = 0.05$ and Prandtl number $Pr = 151$.

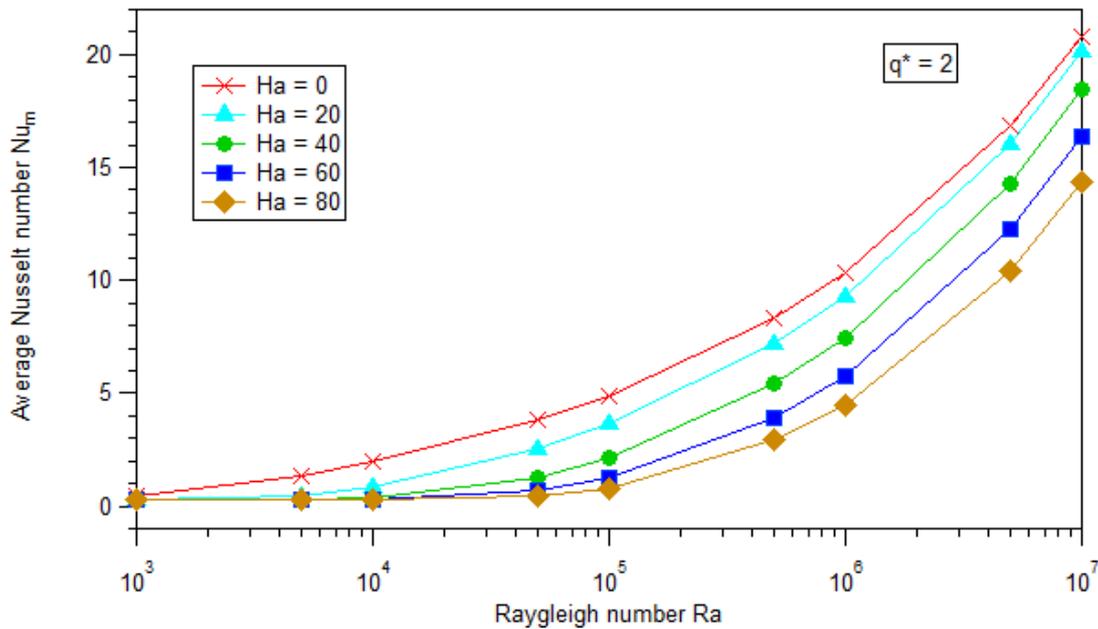


Figure 5. Average Nusselt number according to the Rayleigh number for different Hartmann number at dimensionless heat generation $q^* = 2$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

Figure 6 shows the ratio between average Nusselt number and average Nusselt number without magnetic field $Nu_m/Nu_m(Ha=0)$ according to the Rayleigh number ($Ra = 10^3, 10^4, 10^5, 10^6$ and 10^7) for different Hartmann number ($Ha =$

0, 20, 40, 60 and 80) at dimensionless heat generation $q^* = 2$, a solid volume fraction $\phi = 0.05$ and Prandtl number $Pr = 151$.

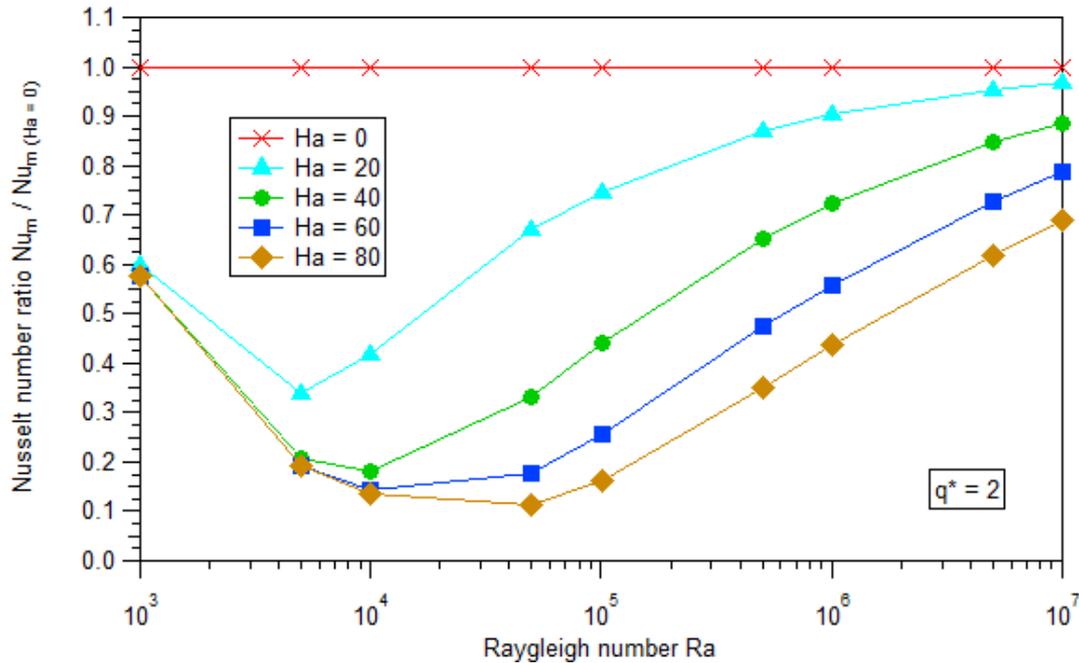


Figure 6. Nusselt number ratio according to the Rayleigh number for different Hartmann number at dimensionless heat generation $q^* = 2$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From these figures, we see that for constant Hartmann number, the Nusselt number increases with the increasing of the Rayleigh number. Also, for constant Rayleigh number, the Nusselt number decreases when the Hartmann number increases.

In addition, due to effect of magnetic field, there are two different phases when the Rayleigh number increases. The first, the Nusselt number ratio decreases to minimum value at critical Rayleigh number depending on the Hartmann number. The second, due to the strong effect of the natural convection

with respect to the effect of magnetic field, this ratio increases.

4.3. Effect of Heat Generation or Absorption and Rayleigh Number on the Isothermal and Streamline Contours

Figure 7 and 8 show the isotherms for dimensionless heat absorption, dimensionless heat generation and different

Rayleigh at Hartmann number $Ha = 40$, a solid volume fraction $\phi = 0.05$ and Prandlt number $Pr = 151$.

Figure 9 and 10 show the streamlines for dimensionless heat absorption, dimensionless heat generation and different Rayleigh at Hartmann number $Ha = 40$ and a solid volume fraction $\phi = 0.05$ and Prandlt number $Pr = 151$.

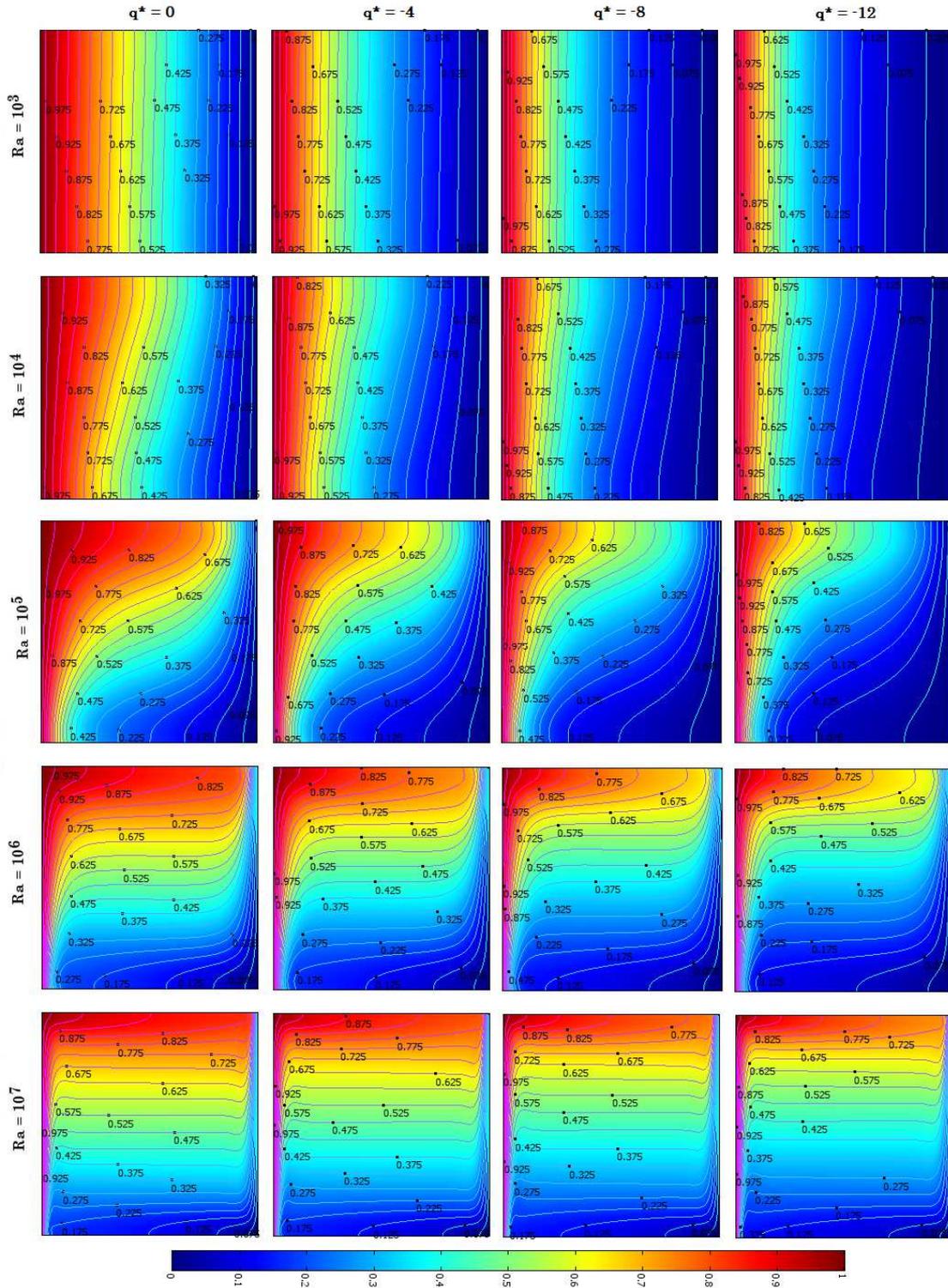


Figure 7. Isotherms for $q^* \leq 0$ and different Rayleigh at Hartmann number $Ha = 40$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

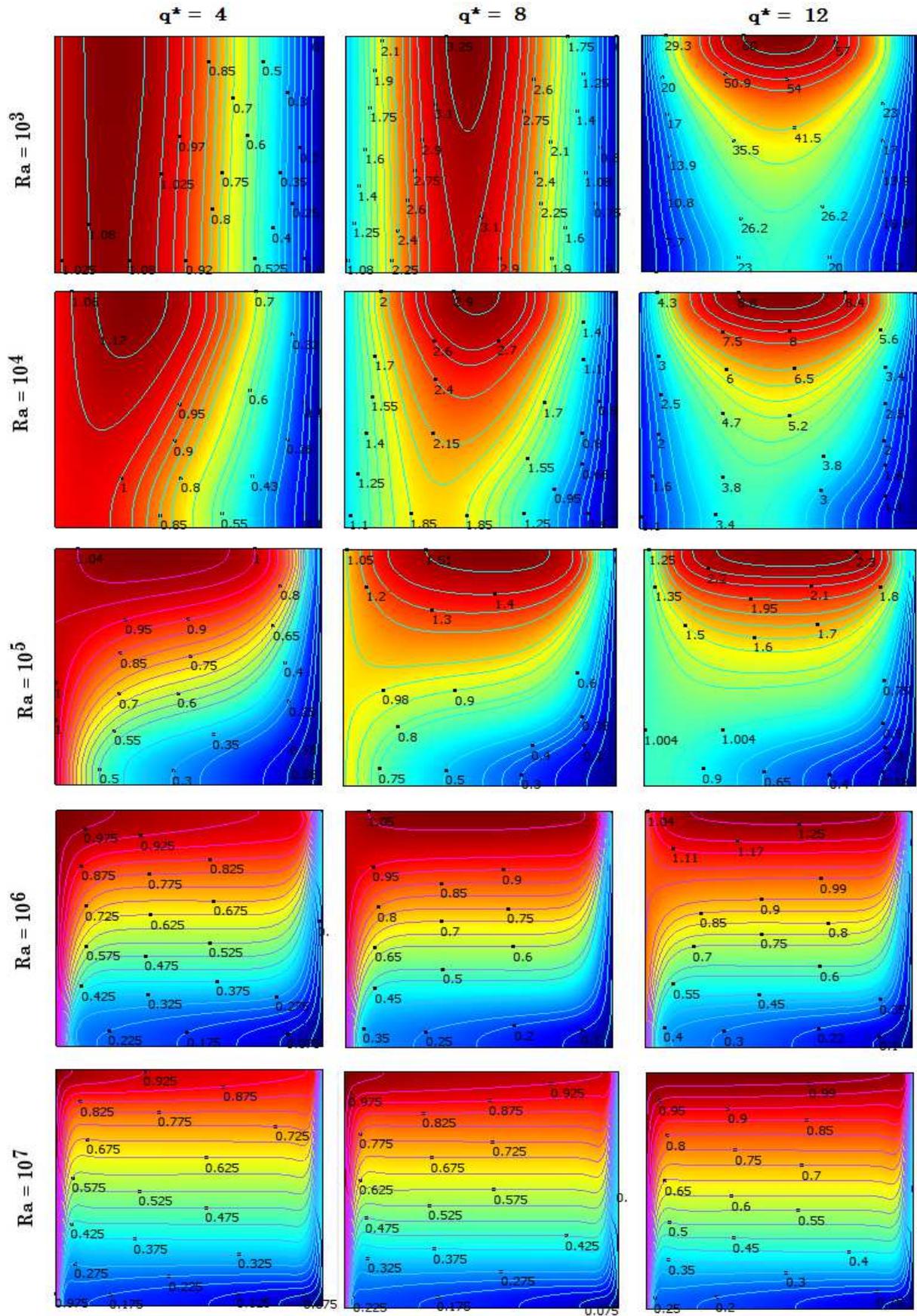


Figure 8. Isotherms for $q^* > 0$ and different Rayleigh at Hartmann number $Ha = 40$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

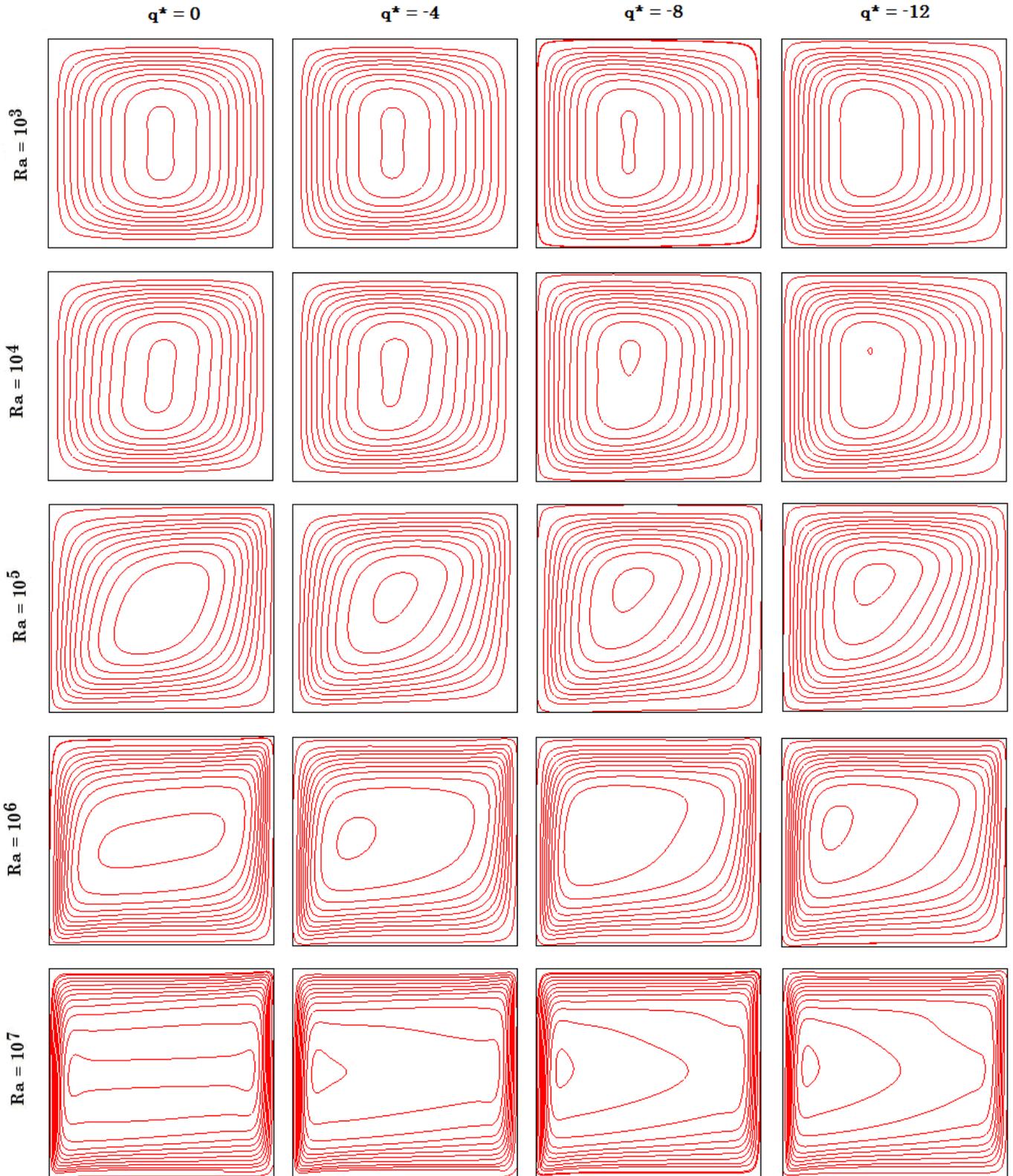


Figure 9. Streamlines for $q^* \leq 0$ and different Rayleigh at Hartmann number $Ha = 40$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

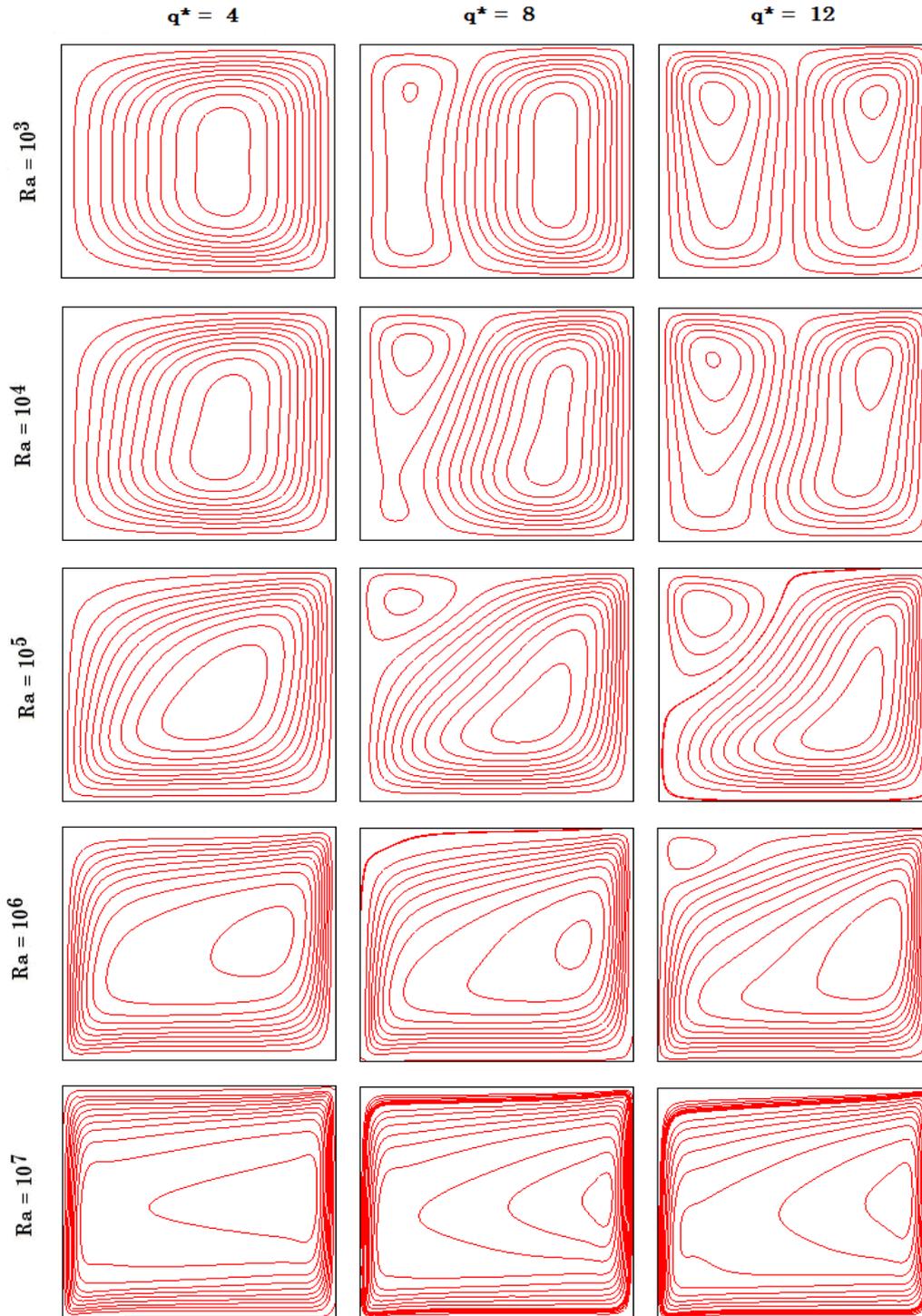


Figure 10. Streamlines for $q^* > 0$ and different Rayleigh at Hartmann number $Ha = 40$ and a solid volume fraction $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From these figures, we see that the conduction regime is dominant with vertical isotherms for the low Rayleigh number $Ra = 10^3$. Also, we observed a strong clockwise circulation in the cavity for a heat sink condition ($q^* = -12$). But, the strength is reduced gradually when the dimensionless heat increases until high heat generation condition ($q^* = 12$). In the latter case, we notice a higher temperature of nanofluid in the cavity which reduces the rate of heat transfer, Nusselt number and the temperature gradient

near the hot wall. Indeed, the two vertical walls temperature is lower than the nanofluids temperature. Therefore, the cold and hot walls receive heat from the nanofluids and a negative Nusselt number at hot wall is found.

In addition, the flow cell becomes stronger when Rayleigh number increases to 10^5 which the convection mode begins to appear. But, this flow cell is reduced when the dimensionless heat q^* increases.

The convection mode begins to appear when Rayleigh

number increased to 10^5 . The flow cell becomes stronger but reduces as q^* increases.

The convection is become dominant when the Rayleigh number reaches the value of 10^7 . The circulating cell is become very strong with a neglected or small for heat generation.

The temperature gradient near the cavity walls without heat sink condition is lower than that with heat sink. Indeed, we observed, in heat sink condition, that the cavity temperature is lower than the temperature in the cavity without heat sink.

4.4. Effect of Heat Generation or Absorption and Hartmann Number on the Nusselt Number

4.4.1. Effect of Heat Generation or Absorption at Constant Hartmann Number

Figure 11 shows the average Nusselt number according to the Rayleigh number for different dimensionless both heat absorption and heat generation at Hartmann number $Ha = 40$ and a solid volume fraction of $\phi = 0.05$ and Prandtl number $Pr = 151$.

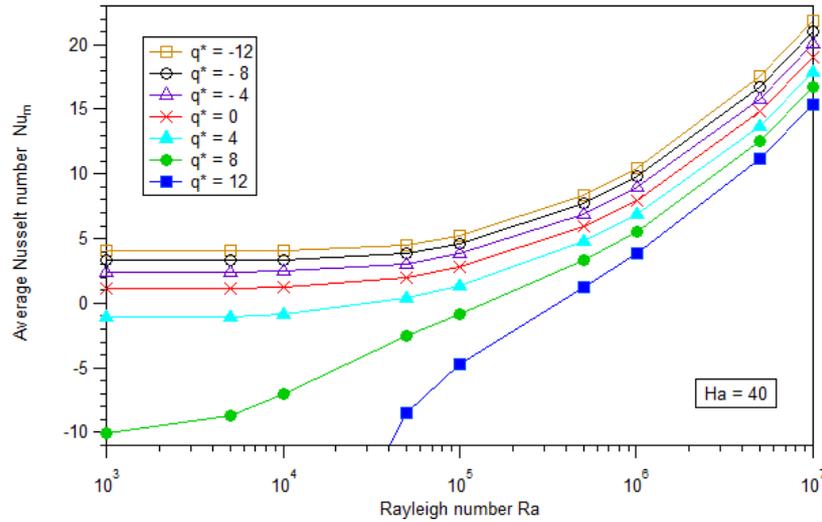


Figure 11. Average Nusselt number according to the Rayleigh number for different dimensionless both heat absorption and heat generation at Hartmann number $Ha = 40$ and a solid volume fraction of $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From this figure, we see that, for constant Rayleigh number, when the dimensionless heat q^* decreases the Nusselt number increases and conversely. Also, for small Rayleigh number, the effect of heat generation ($q^* > 0$) overcomes the effect of natural convection lead to small heat transfer from hot to cold wall or can causes reverse heat transfer from nanofluid to both vertical wall which has a negative Nusselt number. In other word, the effect of both heat generation or absorption is neglected when Rayleigh

number increases.

To study well the effect of heat generation or absorption on the Nusselt number, the Nusselt number ratio (the ratio between Nusselt number and Nusselt number without heat generation $Nu_m/Nu_m(q^*=0)$) according to the Rayleigh number for different dimensionless both heat absorption and heat generation at Hartmann number $Ha = 40$ and a solid volume fraction of $\phi = 0.05$ is shown in Figure 12.

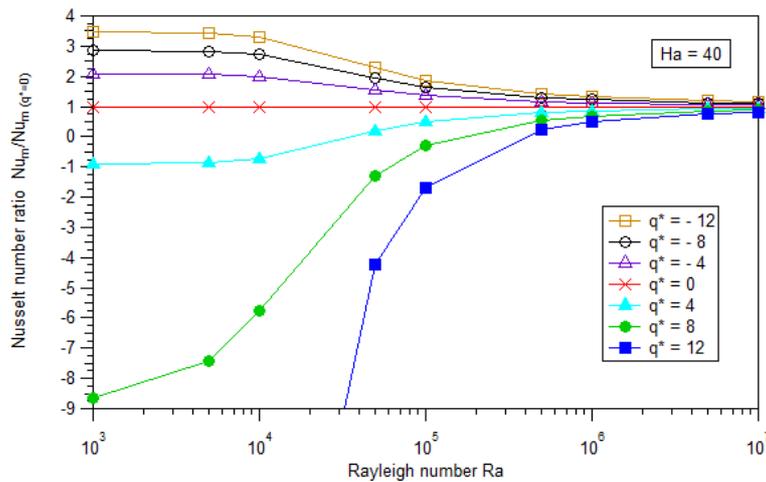


Figure 12. Nusselt number ratio according to the Rayleigh number for different dimensionless both heat absorption and heat generation at Hartmann number $Ha = 40$ and a solid volume fraction of $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From this figure, we see that, for all Rayleigh number and the heat generation, the Nusselt number ratio is less than one but it is greater than unity for the case of heat absorption. Also, for small Rayleigh $Ra = 10^3$, the heat generation or absorption has a major effect on the Nusselt number ratio. This effect decreases with increasing the number of Rayleigh until his neglect at $Ra = 10^7$.

4.4.2. Effect of Heat Generation or Absorption at Constant Rayleigh Number

Figure 13 shows the average Nusselt number according to the Hartmann number for different dimensionless both heat absorption and heat generation at a solid volume fraction of $\phi = 0.05$, Rayleigh number $Ra = 10^5$ and Prandlt number $Pr = 151$.

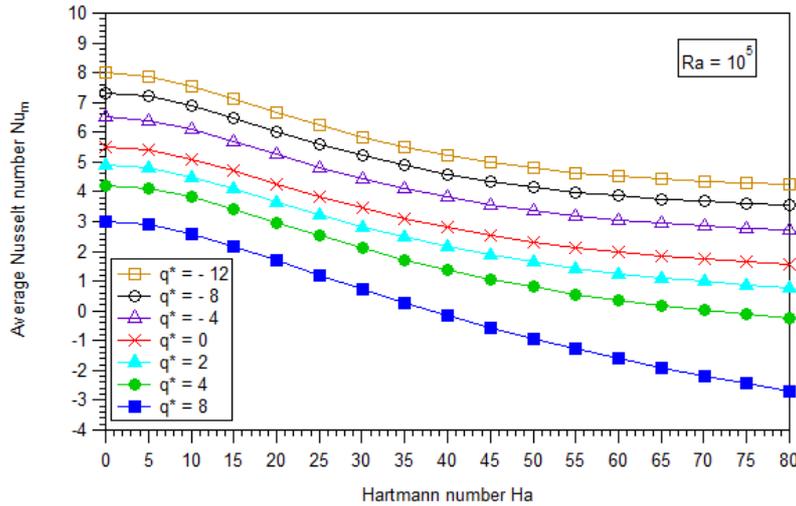


Figure 13. Average Nusselt number according to the Hartmann number for different dimensionless both heat absorption and heat generation at Rayleigh number $Ra = 10^5$ and solid volume fraction of $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From this figure, we always see that the Nusselt number for the case of heat source is lower than the case of heat of absorption. Also, for both heat source or heat absorption, the Nusselt number decreases by increasing the number of Hartmann. Therefore, the increasing of the magnetic field reduces the Nusselt number. This reduction is important for the case of heat generation compared to the case of heat absorption. On the other hand, the effect of magnetic field is

small in the case of heat absorption but it is greater in the case of heat generation.

Figure 14 shows the Nusselt number ratio according to the Hartmann number for different dimensionless both heat absorption and heat generation at a solid volume fraction of $\phi = 0.05$, Rayleigh number $Ra = 10^5$ and Prandlt number $Pr = 151$.

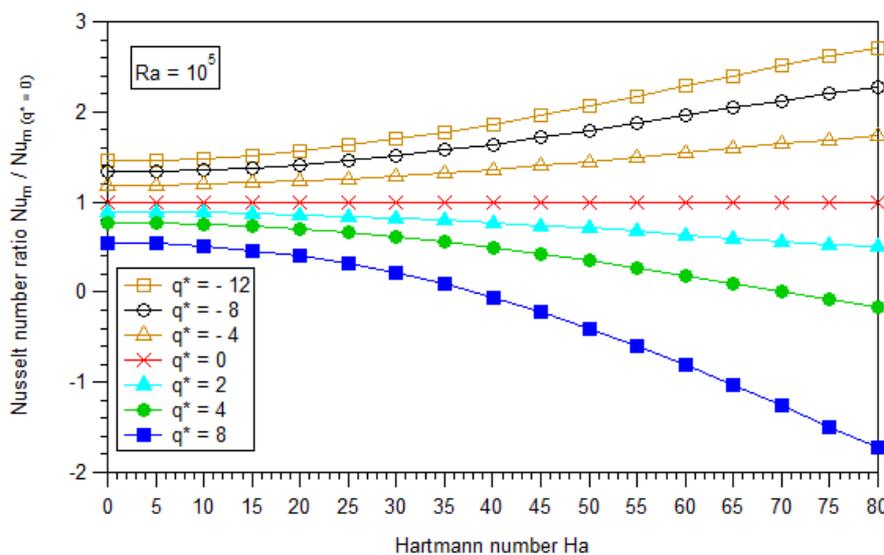


Figure 14. Nusselt number ratio according to the Hartmann number for different dimensionless both heat absorption and heat generation at Rayleigh number $Ra = 10^5$ and solid volume fraction of $\phi = 0.05$ (Ethylene Glycol - Cu nanofluid).

From this figure, we see that, for the situation of the heat absorption, the Nusselt number ratio is greater than unity and

increases as the Hartmann number increases. Therefore, the increasing of the magnetic field raises the Nusselt number ratio. Then, for the case of the heat generation, the Nusselt number ratio is less than unity and decreases as the Hartmann

number increases. Therefore, the increasing of the magnetic field reduces the Nusselt number ratio.

4.5. Comparison with Other Nanofluids

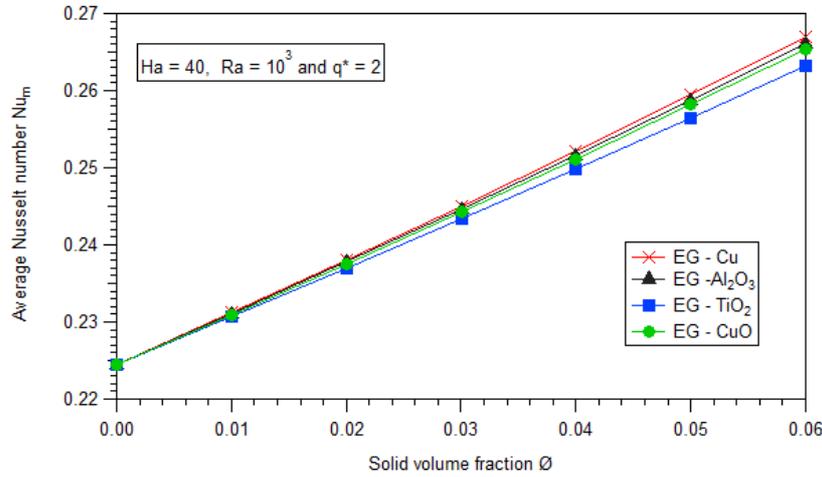


Figure 15. Variation of the average Nusselt number Nu_m according to the solid volume fraction ϕ at different nanofluids for $Ra = 10^3$, $Ha = 40$ and $q^* = 2$.

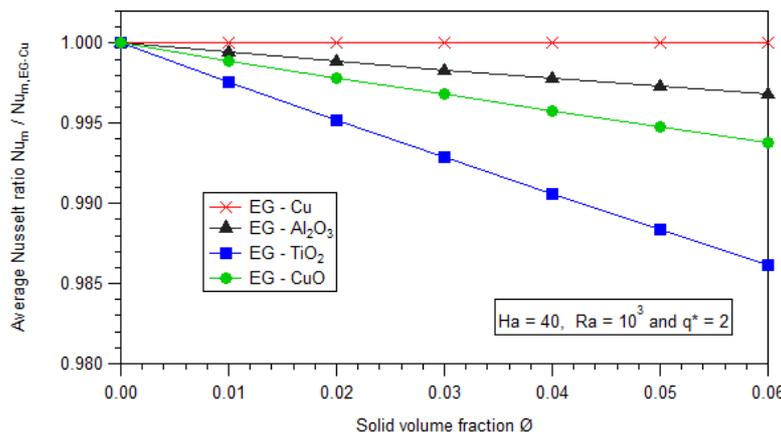


Figure 16. Variation of the average Nusselt ratio $Nu_m / Nu_{m,Ethylene Glycol - cu}$ according to the solid volume fraction ϕ at different nanofluids for $Ra = 10^3$, $Ha = 40$ and $q^* = 2$.

Figure 15 and 16 show, respectively, the variation of the average Nusselt number Nu_m and the average Nusselt number ratio ($Nu_m / Nu_{m,Ethylene Glycol - cu}$) according to the solid volume fraction ($\phi = 0.01; 0.02; 0.03; 0.04; 0.05$ and 0.06) at different nanoparticles (Cu, Al_2O_3 , CuO and TiO_2) for the Hartmann number $Ha = 40$, the Rayleigh number $Ra = 10^3$ and the dimensionless heat generation $q^* = 2$. The choice of the value of Rayleigh number is justified by the dominance of conduction phenomenon for this value. The thermophysical properties of the nanoparticles are given in Table 1. From these figures, we show that the heat transfer depends strongly on the nano thermal conductivity, so Ethylene Glycol-Cu nanofluid enhances the heat transfer compared with Ethylene Glycol - Al_2O_3 , Ethylene Glycol - TiO_2 and Ethylene Glycol - CuO.

5. Conclusion

Natural convection of an electrically conducting fluid in a

square enclosure under an externally imposed uniform horizontal magnetic field and uniform heat generation or absorption is investigated numerically. The two opposing side walls are differentially heated with a different temperatures, whilst the floor and ceiling are thermally insulated. The coupled mass, momentum and energy equations associated are solved the commercial simulation software COMSOL Multiphysics.

The main results obtained of the numerical analysis are as follows :

- For the Hartmann number ($Ha = 40$) and the low Rayleigh number ($Ra = 10^3$ where the conduction regime is dominant), there is a strong clockwise circulation in the cavity for a heat sink condition ($q^* = -12$). But, the strength is reduced gradually when the dimensionless heat increases until high heat generation condition ($q^* = 12$). Therefore, a higher temperature of nanofluid in the cavity which reduces the rate of heat transfer, Nusselt number and the temperature gradient

near the hot wall

- For the Hartmann number ($Ha = 40$) and the higher Rayleigh number ($Ra = 10^7$ where the convection regime is dominant), the cavity temperature is lower than the temperature in the cavity without heat sink (in heat sink condition).
- For the Hartmann number ($Ha = 40$) and constant solid volume fraction ($\phi = 0.05$): for small Rayleigh number, the effect of heat generation ($q^* > 0$) overcomes the effect of natural convection lead to small heat transfer from hot to cold wall or can causes reverse heat transfer from nanofluid to both vertical wall which has a negative Nusselt number. In other word, the effect of both heat generation or absorption is neglected when Rayleigh number increases.
- Compared with other nanoparticles, Ethylene Glycol-Cu nanofluid enhances the heat transfer which depends strongly on the nano thermal conductivity for constant Hartmann number ($Ha = 40$) and constant Rayleigh number ($Ra = 10^3$).

Nomenclature

B_0 : magnetic field strength
 C_p : specific heat
 g : gravitational acceleration
 h : convection heat transfer coefficient
 Ha : Hartmann number
 k : thermal conductivity
 L : enclosure length
 Nu_y : local Nusselt number on the left hot wall
 Nu_m : average Nusselt number
 p : fluid pressure
 p^* : dimensionless pressure
 Pr : Prandtl number
 q : heat generation/absorption
 q^* : dimensionless heat generation/absorption
 Ra : Rayleigh number
 T : temperature, K
 T^* : dimensionless temperature
 u, v : velocity components in x, y directions
 u^*, v^* : dimensionless velocity components
 x, y : cartesian coordinates
 x^*, y^* : dimensionless coordinates

Greek symbols

ρ : thermal diffusivity
 β : thermal expansion coefficient
 ϕ : solid volume fraction
 μ : dynamic viscosity
 ν : kinematic viscosity
 ρ : density
 σ : electrical conductivity
 Ψ : stream function

Subscripts

c : cold wall
 f : fluid
 h : hot wall
 nf : nanofluid
 p : nanoparticle

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